



Formation of Structure in the Universe

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Date	Торіс	Date	Торіс	Date	Торіс
23/01/2023	Introduction	13/02/2023	From gas cloud to collapsed object	06/03/2023	Gravitational instabilities in the cosmological context
27/01/2023	Physical process in baryonic gas (part 1)	17/02/2023	Galaxies and star-formation on galactic scales (part 1)	10/03/2023	Hierarchical structure formation (part 1)
30/01/2023	Physical process in baryonic gas (part 2)	20/01/2023	Galaxies and star-formation on galactic scales (part 2)	13/03/2023	Hierarchical structure formation (part 2)
03/02/2023	Gravitational stability and instability (part 1)	24/02/2023	Galaxies and star-formation on galactic scales (part 3)	17/03/2023	Galaxy formation and evolution
06/02/2023	Gravitational stability and instability (part 2)	27/02/2023	Feedback processes in star formation	April 2023	Exam
10/02/2023	Gravitational collapse	03/03/2023	Galaxies interaction and triggering star-formation		

Supervison

3 groups :

- 6-8 students/group
- please sign in (email circulating soon)
- 3 sessions

Group ID	Session 1		Session 2		Session 3	
	Date	Room	Date	Room	Date	Room
1						
2						
3						

Introduction

Chapter 1



t=0s : *the Big-Bang* **t=10**⁻³⁶ **to t=10**⁻³²**s** : *inflation*

- emission of gravitational waves

- emission of density waves

Guzzetti et al. 2016, arXiv: 1605.01615







Resolution of current instrumentation : 1/1000 of the size of a proton.





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- Universe mainly composed of quarks, leptons and photons t=10⁻⁶s : formation of protons and neutrons







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Deuterium

Lithium



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- Universe mainly composed of quarks, leptons and photons **t=10**⁻⁶**s** : formation of protons and neutrons, and then formation of nuclei of deuterium, helium and lithium.

t=3mn : the Universe is mainly composed of radiation, baryonic matter, dark matter and dark energy.

- Electrons and nuclei are not bounded yet.



- Over-dense regions grow from the initial perturbation
- As the Universe is expanding, the Universe's temperature is decreasing.
- The Universe's density is also decreasing as the Universe is expanding



 Free electrons and protons start to be bounded but as soon as they are hit by a photon they become again unbounded :

$e + p \Leftrightarrow H + \gamma$

- When the Universe's density is sufficiently low to avoid interaction between photons and particles (*matter and radiation are decoupled*):
 - The first atoms are formed (Hydrogen, Helium, Lithium)
 - Photons can escape, and form the first emitted radiation in the Universe : the Cosmic Microwave Background

This epoch is called the recombination phase





- European Space Agency (ESA) satellite
- Launched by an Ariane 5 rocket in May 2009
- Diameter :1.5m
- Objectives :
 - Map the Cosmic Microwave Background
 - Measure the cosmological parameters
 - Study galaxy clusters

- After the recombination phase, the Universe enters the Dark Ages
- Overdense regions continue to grow
- Their density becomes sufficiently large that their gravitational field is dominated by their own mass
- Their evolution is now driven by their own gravity (self-gravitating objects) and not the evolution of the Universe (Hubble flow)
- At the center of these overdense regions, the gas cools and leads to the formation of the first stars : this epoch is called **Cosmic Dawn**



- The first stars are surrounded by neutral hydrogen
- The first stars emit UV photons^{*} which will ionise the neutral hydrogen, creating bubbles of ionised hydrogen : this is the epoch of reionisation



* This will be discussed later in this course

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- CMB observations by *Planck* reveals that the hydrogen is fully ionised 1 billion years after the Big-Bang (z=6)
- The study of the first generation of galaxies shows that they grow hierarchically by merging, leading to an evolving distribution of galaxies of different masses.



Robertson et al. (2015), ApJ, 802, 19

^{*} This will be discussed later in this course



The "hunt" for the most distant galaxies between 1950s and 2022



Humason et al. 1956 Minkowski 1960 Spinrad et al. 1975 Spinrad & Smith 1976 Smith et al. 1979 Spinrad 1982 Lilly 1988 Chambers et al. 1990 Lacy et al. 1994 Petitjean et al. 1996 Franz et al. 1997 Day et al. 1998 Hu et al. 1999, 2002 Pelló et al. 2004 lye et al. 2006 Fontana et al. 2010 Vanzella et al. 2011 Ono et al. 2012 Shibuya et al. 2012 Finkelstein et al. 2013 Oesch et al 2014 Zitrin et al. 2015 Oesch et al. 2016 Harikane et al. 2022

















GLASS-z13 (Naidu et al. 2022)



SMACS-z16a (Atek et al. 2022)

Within a week, 11 papers have been submitted using the first dataset from the JWST to search for the first galaxies. z=17



CEERS-DSFG-1 (Finkelstein et al. 2022)

Summary of the formation of structures in the Universe



processes responsible for the cooling of the gas, the heating of the gas and the formation of protostars.

Toolbox for this course



Equation of continuity



Poisson's equation

Differential equation

 $\Delta \Phi = \nabla^2 \Phi = f$

For the gravitational potential we can write :

 $\nabla^2 \Phi_g = 4\pi g \rho(\vec{r})$

Equation of state for an ideal gas



Physical processes in baryonic gas

Chapter 2 (part 2)

Some definitions

The amount of energy (dE) passing through a surface should be proportional to the size of the surface (dA)and to the duration of the exposition (dt).

It is usually defined as :

 $dE = dF \times dA \times dt$

The energy flux dF is measured in erg s⁻¹ cm⁻²

A source of radiation is isotropic if it emits energy equally in all directions

By conservation of Energy, $dE(r_1) = dE(r_2)$ then : $F(r_1)4\pi r_1^2 = F(r_2)4\pi r_2^2$



$$F(r) = \frac{F(r_1)r_1^2}{r^2}$$

Some definitions

Considering a surface normal to the direction of a given ray, and considering all the rays whose direction is within a solid angle $d\Omega$, then the energy passing through the element dA is :

 $dE = I_{\nu} dA dt d\nu d\Omega$ where I_{ν} is the <u>specific intensity</u> (or brightness)





If the surface is not perpendicular to the rays but has different orientation, then $dE = I_{\nu} dA dt d\nu \cos \theta d\Omega$ $= dF_{\nu} dA dt d\nu$

<u>N.B.</u>: If the radiation is isotropic, then $\int dF_{\nu} = 0$

Some definitions

The <u>specific energy density</u> is the energy per unit volume per unit frequency range :

 $dE = u_{\nu}(\Omega) \, dV \, d\Omega \, d\nu$

If we consider this cylinder, then dV = dA c dT then $dE = u_{\nu}(\Omega) dA c dt d\Omega d\nu$ but within dt all radiation will pass out of the cylinder, then :

$$dE = I_{\nu} dA \, d\Omega \, dt \, d\nu$$

The specific energy density is defined as :

$$u_{\nu}(\Omega) = \frac{I_{\nu}}{c}$$

The mean density is defined as :

$$J_{\nu} = \frac{1}{4\pi} \int I_{\nu} d\Omega$$



Integrating the specific energy density over solid angles :

$$u_{\nu} = \int u_{\nu}(\Omega) d\Omega = \frac{1}{c} \int I_{\nu} d\Omega$$
 or $u_{\nu} = \frac{4\pi}{c} J_{\nu}$

The total radiation density is given by
$$u = \int u_{\nu} d\nu = \frac{4\pi}{c} \int J_{\nu} d\nu$$

Radiative transfer

We use radiative transfer each time a radiation is passing through a matter and add (absorption) or subtract (emission) energy.

I- EMISSION

The <u>spontaneous emission coefficient</u> is defined as the energy emitted per unit time per unit solid angle per unit volume, such as : $dE = j \ dV \ d\Omega \ dt$

If the emission is monochromatic (e.g., an emission line), we can define a monochromatic emission coefficient :

 $dE = j_{\nu} dV \, d\Omega \, dt \, d\nu$

If the emission is isotropic, then :

$$j_{\nu} = \frac{1}{4\pi} P_1$$

where P_{ν} is the radiated power

We can also define <u>the emissivity</u> as the energy emitted per unit frequency per unit time per unit mass, and rewrite the transmitted energy in an isotropic emission as :

$$dE = \epsilon_{\nu} \rho \, dV \, dt \, d\nu \, \frac{d\Omega}{4\pi}$$

For an isotropic emission, the relation between the emission coefficient and the emissivity is given by :

$$j_{\nu} = \frac{\epsilon_{\nu}\rho}{4\pi}$$

Considering a beam of cross-section dA traveling through a volume $dV = ds \times dA$, then the energy added by spontaneous emission is : $dI_{v} = i_{v} ds$

(Remember that $dE = I_{\nu} dA dt d\nu d\Omega$)

Radiative transfer

We use radiative transfer each time a radiation is passing through a matter and add (absorption) or subtract (emission) energy.

II- ABSORPTION

The <u>absorption coefficient</u> is defined as the loss of intensity in a beam as it travels a distance ds: $dI_{\nu} = -\alpha_{\nu} I_{\nu} ds$

The absorption depends on the density of absorbers along the travel of a beam.

If we assume a random distribution of absorbers, each of them with a cross section σ_{ν} and a density per unit volume n, then the effect of these absorbers on a radiation passing through dA within a solid angle $d\Omega$ is :

 $dE = -dI_{\nu} \, dA \, dt \, d\Omega \, d\nu$

or : $dE = I_{\nu}(n \, dA \, ds \, \sigma_{\nu}) d\Omega \, dt \, d\nu$

Therefore :

$$dI_{\nu} = -n \sigma_{\nu} I_{\nu} ds$$

We can rewrite the absorption coefficient such as : $\alpha_{\nu} = n \: \sigma_{\nu}$

Usually, α_{ν} is defined with the opacity (also known as the mass absorption coefficient) such as :

 $\alpha_{\nu} = \rho \kappa_{\nu}$